Lecture 5: Stochastic Runge–Kutta Methods

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1 Introduction

- 2 Runge–Kutta methods for ODEs
- Strong stochastic Runge–Kutta methods
- Weak stochastic Runge–Kutta methods



Overview of this lecture

Runge–Kutta methods for ODEs

- Taylor series.
- General Runge-Kutta schemes.
- Explicit and implicit schemes.
- Strong stochastic Runge–Kutta methods
 - Itô–Taylor series.
 - A family of strong order 1.0 schemes.
 - The iterated Itô integrals.
- Weak stochastic Runge–Kutta methods
 - A family of weak order 2.0 schemes.
 - Approximating the iterated Itô integrals.

- A family of iterative methods for solving differential equations.
- Based on Taylor series (see the previous lecture), ...
- ... but are derivative-free.
- Plug-and-play methods that only require specification of the differential equation (at least ideally).
- There are other methods as well (not considered here):
 - Multistep methods (e.g. Adams methods)
 - Multiderivative methods
 - Higher-order methods (e.g. Nyström method)
 - Tailored methods (for specific problems)

• Consider a first-order non-linear ODE

$$\frac{\mathrm{d}\mathbf{x}(t)}{\mathrm{d}t} = \mathbf{f}(\mathbf{x}(t), t), \qquad \mathbf{x}(t_0) = \text{given},$$

- The simplest Runge–Kutta method is the (forward) Euler scheme.
- It is based on sequential linearization of the ODE system:

$$\hat{\mathbf{x}}(t_{k+1}) = \hat{\mathbf{x}}(t_k) + \mathbf{f}(\hat{\mathbf{x}}(t_k), t_k) \Delta t.$$

- Easy to understand and implement.
- The global error of the method depends linearly on the step size Δt .

Taylor series [1/2]

• The ODE system can be integrated to give

$$\mathbf{x}(t) = \mathbf{x}(t_0) + \int_{t_0}^t \mathbf{f}(\mathbf{x}(\tau), \tau) \, \mathrm{d}\tau.$$

 Recall from the previous lecture that we used a Taylor series expansion for the solution of the ODE

$$\begin{aligned} \mathbf{x}(t) &= \mathbf{x}(t_0) + \mathbf{f}(\mathbf{x}(t_0), t_0) (t - t_0) \\ &+ \frac{1}{2!} \, \mathcal{L} \, \mathbf{f}(\mathbf{x}(t_0), t_0) (t - t_0)^2 \\ &+ \frac{1}{3!} \, \mathcal{L}^2 \, \mathbf{f}(\mathbf{x}(t_0), t_0) (t - t_0)^3 + \dots \end{aligned}$$

We used the linear operator

$$\mathcal{L}(\bullet) = \frac{\partial}{\partial t}(\bullet) + \sum_{i} f_{i} \frac{\partial}{\partial x_{i}}(\bullet)$$

.

In other words, the series expansion is equal to

$$\begin{aligned} \mathbf{x}(t) &= \mathbf{x}(t_0) + \mathbf{f}(\mathbf{x}(t_0), t_0) \left(t - t_0\right) \\ &+ \frac{1}{2!} \left\{ \frac{\partial}{\partial t} \mathbf{f}(\mathbf{x}(t_0), t_0) + \sum_i f_i(\mathbf{x}(t_0), t_0) \frac{\partial}{\partial x_i} \mathbf{f}(\mathbf{x}(t_0), t_0) \right\} \left(t - t_0\right)^2 \\ &+ \frac{1}{3!} \left\{ \frac{\partial [\mathcal{L} \mathbf{f}(\mathbf{x}(t_0), t_0)]}{\partial t} + \sum_i f_i(\mathbf{x}(t_0), t_0) \frac{\partial [\mathcal{L} \mathbf{f}(\mathbf{x}(t_0), t_0)]}{\partial x_i} \right\} \left(t - t_0\right)^3 \\ &+ \dots \end{aligned}$$

 If we were only to consider the terms up to ∆t, we would recover the Euler method.

Derivation of a higher-order method [1/4]

- However, here we wish to get hold of higher-order methods.
- For the sake of simplicity, we now stop at the term $(t t_0)^2 = (\Delta t)^2$.
- We get

$$\mathbf{x}(t_0 + \Delta t) \approx \mathbf{x}(t_0) + \mathbf{f}(\mathbf{x}(t_0), t_0) \Delta t + \frac{1}{2} \left\{ \frac{\partial}{\partial t} \mathbf{f}(\mathbf{x}(t_0), t_0) + \sum_i f_i(\mathbf{x}(t_0), t_0) \frac{\partial}{\partial x_i} \mathbf{f}(\mathbf{x}(t_0), t_0) \right\} (\Delta t)^2.$$

 We aim to get rid of the derivatives and be able to write the expression in terms of the function f(·, ·) evaluated at various points. • We now seek a form with an extra stage:

 $\mathbf{x}(t_0 + \Delta t) \approx \mathbf{x}(t_0) + A \mathbf{f}(\mathbf{x}(t_0), t_0) \Delta t$ $+ B \mathbf{f}(\mathbf{x}(t_0) + C \mathbf{f}(\mathbf{x}(t_0), t_0) \Delta t, t_0 + D \Delta t) \Delta t,$

where A, B, C, and D are unknown.

• In the last term, we can consider the truncated Taylor expansion (linearization) around $\mathbf{f}(\mathbf{x}(t_0), t_0)$ with the chosen increments as follows:

$$\mathbf{f}(\mathbf{x}(t_0) + C \mathbf{f}(\mathbf{x}(t_0), t_0) \Delta t, t_0 + D \Delta t) = \mathbf{f}(\mathbf{x}(t_0), t_0) + C \left(\sum_i f_i(\mathbf{x}(t_0), t_0) \frac{\partial}{\partial x_i} \mathbf{f}(\mathbf{x}(t_0), t_0) \right) \Delta t + D \frac{\partial \mathbf{f}(\mathbf{x}(t_0), t_0)}{\partial t} \Delta t + \cdots$$

• Combining the previous two equations gives:

$$\mathbf{x}(t_0 + \Delta t) \approx \mathbf{x}(t_0) + (\mathbf{A} + \mathbf{B}) \mathbf{f}(\mathbf{x}(t_0), t_0) \Delta t + \mathbf{B} \left[C \sum_i f_i(\mathbf{x}(t_0), t_0) \frac{\partial}{\partial x_i} \mathbf{f}(\mathbf{x}(t_0), t_0) + D \frac{\partial \mathbf{f}(\mathbf{x}(t_0), t_0)}{\partial t} \right] (\Delta t)^2$$

 If we now compare the above equation to the original truncated Taylor expansion, we get the following conditions for our coefficients:

$$A + B = 1$$
, $B = \frac{1}{2}$, $C = 1$, and $D = 1$.

Derivation of a higher-order method [4/4]

What we derived here is a two-stage method (known as Heun's method):

$$\hat{\mathbf{x}}(t_0 + \Delta t) = \mathbf{x}(t_0) + \frac{\Delta t}{2} \{ \mathbf{f}(\tilde{\mathbf{x}}_1, t_0) + \mathbf{f}(\tilde{\mathbf{x}}_2, t_0 + \Delta t) \},\$$

where the supporting values are given by

$$\begin{split} \tilde{\mathbf{x}}_1 &= \mathbf{x}(t_0), \\ \tilde{\mathbf{x}}_2 &= \mathbf{x}(t_0) + \mathbf{f}(\tilde{\mathbf{x}}_1, t_0) \,\Delta t. \end{split}$$

- The method (in practice the finite differences) are determined by the choices we did in truncating the series expansion.
- This method is of order 2.

Algorithm: Runge-Kutta method

Start from $\hat{\mathbf{x}}(t_0) = \mathbf{x}(t_0)$ and divide the integration interval $[t_0, t]$ into n steps $t_0 < t_1 < t_2 < \ldots < t_n = t$ such that $\Delta t = t_{k+1} - t_k$. The integration method is defined by its Butcher tableau:

On each step *k* approximate the solution as follows:

$$\hat{\mathbf{x}}(t_{k+1}) = \hat{\mathbf{x}}(t_k) + \sum_{i=1}^{s} \alpha_i \mathbf{f}(\tilde{\mathbf{x}}_i, \tilde{t}_i) \Delta t,$$

where $\tilde{t}_i = t_k + c_i \Delta t$ and $\tilde{\mathbf{x}}_i = \hat{\mathbf{x}}(t_k) + \sum_{j=1}^s A_{i,j} \mathbf{f}(\tilde{\mathbf{x}}_j, \tilde{t}_j) \Delta t$.

• Ordinary Runge–Kutta methods are commonly expressed in terms of a table called the Butcher tableau:

$$\begin{array}{c|c} c_1 & A_{1,1} \\ c_2 & A_{2,1} & A_{2,2} \\ \vdots & \vdots & \ddots \\ c_s & A_{s,1} & A_{s,2} & \dots & A_{s,s} \\ \hline & \alpha_1 & \alpha_2 & \dots & \alpha_s \end{array}$$

Example (Forward Euler)

The forward Euler scheme has the Butcher tableau:

which gives the recursion $\hat{\mathbf{x}}(t_{k+1}) = \hat{\mathbf{x}}(t_k) + \mathbf{f}(\hat{\mathbf{x}}(t_k), t_k) \Delta t$.

Example (The fourth-order Runge–Kutta method)

The well-known RK4 method in Ch. 1 has the following Butcher tableau:

- We have considered this far so-called explicit schemes.
- Numerical instability, when the solution includes rapidly varying terms (stiff problems).
- Explicit schemes use very small step sizes in order to not diverge from a solution path (computationally demanding).
- In implicit Runge–Kutta methods, the Buther tableau is no longer lower-triangular.
- On every step, a system of algebraic equations has to be solved (computationally demanding, but more stabile).

• The simplest implicit method is the backward Euler scheme.

Example (Backward Euler)

The implicit backward Euler scheme has the Butcher tableau:

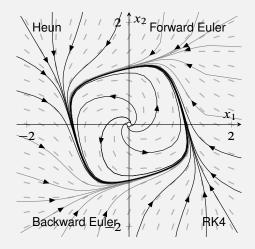
which gives the recursion $\hat{\mathbf{x}}(t_{k+1}) = \hat{\mathbf{x}}(t_k) + \mathbf{f}(\hat{\mathbf{x}}(t_{k+1}), t_k + \Delta t) \Delta t$.

• We study the two-dimensional non-linear ordinary differential equation system

$$\dot{x}_1 = x_1 - x_2 - x_1^3,$$

 $\dot{x}_2 = x_1 + x_2 - x_2^3,$

Example [2/2]



Strong stochastic Runge–Kutta: Basic principles

- A family of iterative methods for solving stochastic differential equations.
- Based on Itô-Taylor series (see the previous lecture), ...
- ... but are derivative-free.
- Plug-and-play methods that only requires specification of the drift and diffusion function of the SDE (at least ideally).
- We divide the methods into strong and weak methods (as we did for the Itô–Taylor series approximations).
- A word of warning: Stochastic Runge–Kutta methods are not as easy to grasp as the ordinary ones.

Itô-Taylor series [1/1]

Recall the following multi-dimensional SDE formulation

 $d\mathbf{x} = \mathbf{f}(\mathbf{x}(t), t) dt + \mathbf{L}(\mathbf{x}(t), t) d\beta, \qquad \mathbf{x}(t_0) \sim p(\mathbf{x}(t_0)),$

where the drift is defined by $\mathbf{f} : \mathbb{R}^d \times \mathbb{R} \to \mathbb{R}^d$ and the diffusion coefficients by $\mathbf{L} : \mathbb{R}^d \times \mathbb{R} \to \mathbb{R}^d \times \mathbb{R}^m$.

The driving noise process

$$\beta(t) = (\beta^{(1)}(t), \beta^{(2)}(t), \dots, \beta^{(m)}(t))$$

is an *m*-dimensional standard Brownian motion.

In integral form the equation can be expressed as

$$\mathbf{x}(t) = \mathbf{x}(t_0) + \int_{t_0}^t \mathbf{f}(\mathbf{x}(\tau), \tau) \, \mathrm{d}\tau + \int_{t_0}^t \mathbf{L}(\mathbf{x}(\tau), \tau) \, \mathrm{d}\beta(\tau).$$

Itô-Taylor series [1/2]

3

 Applying the Itô formula to the terms f(x(t), t) and L(x(t), t) and collecting the terms gives an Itô–Taylor series expansion of the solution (see the previous lecture):

$$\begin{split} \mathbf{x}(t) &= \mathbf{x}(t_0) + \mathbf{f}(\mathbf{x}(t_0), t_0) \left(t - t_0\right) + \mathbf{L}(\mathbf{x}(t_0), t_0) \left(\beta(t) - \beta(t_0)\right) \\ &+ \int_{t_0}^t \int_{t_0}^\tau \mathcal{L}_t \mathbf{f}(\mathbf{x}(\tau), \tau) \, \mathrm{d}\tau \, \mathrm{d}\tau \\ &+ \sum_i \int_{t_0}^t \int_{t_0}^\tau \mathcal{L}_{\beta,i} \mathbf{f}(\mathbf{x}(\tau), \tau) \, \mathrm{d}\beta^{(i)}(\tau) \, \mathrm{d}\tau \\ &+ \int_{t_0}^t \int_{t_0}^\tau \mathcal{L}_t \mathbf{L}(\mathbf{x}(\tau), \tau) \, \mathrm{d}\tau \, \mathrm{d}\beta(\tau) \\ &+ \sum_i \int_{t_0}^t \int_{t_0}^\tau \mathcal{L}_{\beta,i} \mathbf{L}(\mathbf{x}(\tau), \tau) \, \mathrm{d}\beta^{(i)}(\tau) \, \mathrm{d}\beta(\tau). \end{split}$$

- The first row in the equation is just the Euler–Maruyama scheme.
- Similarly as we did in for the ordinary RK methods, we can consider truncated series expansions of various degrees for each of these terms.
- The extra terms involving the iterated and cross-term Itô integrals complicate the formulation.
- To present a family of actual numerical methods, we consider the following family of strong order 1.0 methods due to Andreas Rößler...

Start from x̂(t₀) = x(t₀) and divide the integration interval [t₀, t] into *n* steps t₀ < t₁ < t₂ < ... < t_n = t such that Δt = t_{k+1} - t_k. The integration method is characterized by its *extended Butcher tableau*:

c ⁽⁰⁾	A ⁽⁰⁾	B ⁽⁰⁾	
c ⁽¹⁾	A ⁽¹⁾	B ⁽¹⁾	
	$lpha^{T}$	$\left[oldsymbol{\gamma}^{(1)} ight]^{T}$	$\left[\gamma^{(2)} ight]^{T}$

• On each step *k* approximate the solution trajectory as follows:

$$\hat{\mathbf{x}}(t_{k+1}) = \hat{\mathbf{x}}(t_k) + \sum_{i=1}^{s} \alpha_i \mathbf{f}(\tilde{\mathbf{x}}_i^{(0)}, t_k + c_i^{(0)} \Delta t) \Delta t$$
$$+ \sum_{i=1}^{s} \sum_{n=1}^{m} (\gamma_i^{(1)} \Delta \beta_k^{(n)} + \gamma_i^{(2)} \sqrt{\Delta t}) \mathbf{L}^n (\tilde{\mathbf{x}}_i^{(n)}, t_k + c_i^{(1)} \Delta t)$$

A class of SRK methods of strong order 1.0 [3/3]

With the supporting values

$$\begin{split} \tilde{\mathbf{x}}_{i}^{(0)} &= \hat{\mathbf{x}}(t_{k}) + \sum_{j=1}^{s} A_{i,j}^{(0)} \, \mathbf{f}(\tilde{\mathbf{x}}_{j}^{(0)}, t_{k} + c_{j}^{(0)} \Delta t) \, \Delta t \\ &+ \sum_{j=1}^{s} \sum_{l=1}^{m} B_{i,j}^{(0)} \, \mathbf{L}^{l}(\tilde{\mathbf{x}}_{j}^{(l)}, t_{k} + c_{j}^{(1)} \Delta t) \, \Delta \beta_{k}^{(l)}, \\ \tilde{\mathbf{x}}_{i}^{(n)} &= \hat{\mathbf{x}}(t_{k}) + \sum_{j=1}^{s} A_{i,j}^{(1)} \, \mathbf{f}(\tilde{\mathbf{x}}_{j}^{(0)}, t_{k} + c_{j}^{(0)} \Delta t) \, \Delta t \\ &+ \sum_{j=1}^{s} \sum_{l=1}^{m} B_{i,j}^{(1)} \, \mathbf{L}^{l}(\tilde{\mathbf{x}}_{j}^{(l)}, t_{k} + c_{j}^{(1)} \Delta t) \, \frac{\Delta \beta_{k}^{(l,n)}}{\sqrt{\Delta t}} \end{split}$$

for i = 1, 2, ..., s and n = 1, 2, ..., m.

• The increments in the algorithm are given by the Itô integrals:

$$\Delta eta_k^{(i)} = \int_{t_k}^{t_{k+1}} \mathrm{d} eta^{(i)}(au) \quad \text{and} \ \Delta eta_k^{(i,j)} = \int_{t_k}^{t_{k+1}} \int_{t_k}^{ au_2} \mathrm{d} eta^{(i)}(au_1) \, \mathrm{d} eta^{(j)}(au_2),$$

• The increments $\Delta \beta_k^{(i)}$ are independent normally distributed random variables

 $\Delta \beta_k^{(i)} \sim \mathsf{N}(0, \Delta t).$

- The iterated stochastic Itô integrals $\Delta \beta_k^{(i,j)}$ are trickier.
- For these methods, when *i* = *j*, the multiple Itô integrals can be rewritten as

$$\Delta \beta_k^{(i,i)} = \frac{1}{2} \left(\left[\Delta \beta_k^{(i)} \right]^2 - \Delta t \right).$$

• Exact simulation from the integrals $\Delta \beta_k^{(i,j)}$, when $i \neq j$, is not possible, but can be approximated.

Example (Euler-Maruyama Butcher tableau)

The Euler-Maruyama method has the extended Butcher tableau:

and as we recall from the previous chapter, it is of strong order 0.5.

Example: A strong order 1.0 method

Example (Strong order 1.0 SRK due to Rößler)

• Consider a stochastic Runge–Kutta method with the following extended Butcher tableau:

0									
1	1			0					
0	0	0		0	0				
0									
1	1			1					
1	1	0		-1	0				
	$\frac{1}{2}$	<u>1</u> 2	0	1	0	0	0	$\frac{1}{2}$	$-\frac{1}{2}$

• The (rather lengthy) algorithm is written out in the lecture notes.

- Higher-order methods by considering more terms in the Itô–Taylor expansion.
- Not very practical in general (heavy and complicated).
- For models with some special structure this might still be feasible:
 - One-dimensional models.
 - Additive noise models.
 - Diagonal noise models.
 - Models with commutative noise.

Example: Duffing van der Pol oscillator [1/4]

• Consider a simplified version of a Duffing van der Pol oscillator

$$\ddot{x} + \dot{x} - (\alpha - x^2) x = x w(t), \quad \alpha \ge 0,$$

driven by multiplicative white noise w(t) with spectral density q.

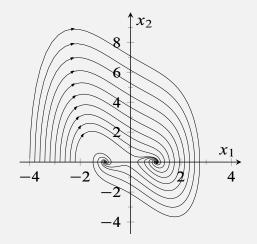
• The corresponding two-dimensional, $\mathbf{x}(t) = (x, \dot{x})$, Itô stochastic differential equation is

$$\begin{pmatrix} \mathrm{d} x_1 \\ \mathrm{d} x_2 \end{pmatrix} = \begin{pmatrix} x_2 \\ x_1(\alpha - x_1^2) - x_2 \end{pmatrix} \mathrm{d} t + \begin{pmatrix} \mathbf{0} \\ x_1 \end{pmatrix} \mathrm{d} \beta,$$

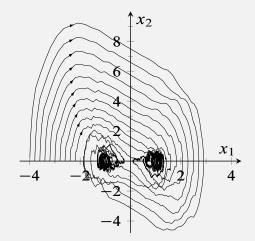
where $\beta(t)$ is a one-dimensional Brownian motion.

• Consider different initializations with q = 0 and $q = 0.5^2$. Use the same realizations of noise for each initialization. Use $\alpha = 1$ and $\Delta t = 2^{-5}$.

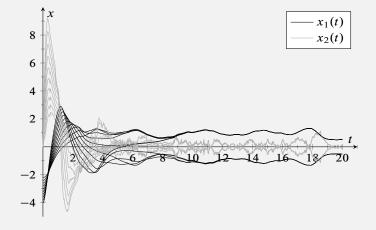
Example: Duffing van der Pol oscillator [2/4]



Example: Duffing van der Pol oscillator [3/4]



Example: Duffing van der Pol oscillator [4/4]



- It is possible to form weak approximations to SDEs, where the interest is not in the solution trajectories, but the distribution of them.
- We can replace weak Itô–Taylor approximations by Runge–Kutta style approximations which avoid the use of derivatives of the drift and diffusion coefficients.
- The reasoning behind the methods is very much the same as for strong SRK methods.
- Here we consider a rather general class of weak order 2.0 methods by Rößler:

A class of SRK methods of weak order 2.0 [1/3]

Start from x̂(t₀) = x(t₀) and divide the integration interval [t₀, t] into *n* steps t₀ < t₁ < t₂ < ... < t_n = t such that Δt = t_{k+1} - t_k. The integration method is characterized by the following extended Butcher tableau:

c ⁽⁰⁾	A ⁽⁰⁾	B ⁽⁰⁾	
c ⁽¹⁾	A ⁽¹⁾	B ⁽¹⁾	
c ⁽²⁾	A ⁽²⁾	B ⁽²⁾	
	$lpha^{T}$	$\left[\gamma^{(1)} ight]^{T}$	$\left[\gamma^{(2)} ight]^{T}$
		$\left[\gamma^{(3)} ight]^{T}$	$\left[\gamma^{(4)} ight]^{T}$

A class of SRK methods of weak order 2.0 [2/3]

• On each step *k* approximate the solution by the following:

$$\hat{\mathbf{x}}(t_{k+1}) = \hat{\mathbf{x}}(t_k) + \sum_{i=1}^{s} \alpha_i \mathbf{f}(\tilde{\mathbf{x}}_i^{(0)}, t_k + \mathbf{c}_i^{(0)} \Delta t) \Delta t$$

$$+ \sum_{i=1}^{s} \sum_{n=1}^{m} \gamma_i^{(1)} \mathbf{L}^n (\tilde{\mathbf{x}}_i^{(n)}, t_k + \mathbf{c}_i^{(1)} \Delta t) \Delta \hat{\beta}_k^{(n)}$$

$$+ \sum_{i=1}^{s} \sum_{n=1}^{m} \gamma_i^{(2)} \mathbf{L}^n (\tilde{\mathbf{x}}_i^{(n)}, t_k + \mathbf{c}_i^{(1)} \Delta t) \frac{\Delta \hat{\beta}_k^{(n,n)}}{\sqrt{\Delta t}}$$

$$+ \sum_{i=1}^{s} \sum_{n=1}^{m} \gamma_i^{(3)} \mathbf{L}^n (\bar{\mathbf{x}}_i^{(n)}, t_k + \mathbf{c}_i^{(2)} \Delta t) \Delta \hat{\beta}_k^{(n)}$$

$$+ \sum_{i=1}^{s} \sum_{n=1}^{m} \gamma_i^{(4)} \mathbf{L}^n (\bar{\mathbf{x}}_i^{(n)}, t_k + \mathbf{c}_i^{(2)} \Delta t) \sqrt{\Delta t},$$

A class of SRK methods of weak order 2.0 [3/3]

• With supporting values

$$\begin{split} \tilde{\mathbf{x}}_{i}^{(0)} &= \hat{\mathbf{x}}(t_{k}) + \sum_{j=1}^{s} A_{i,j}^{(0)} \, \mathbf{f}(\tilde{\mathbf{x}}_{j}^{(0)}, t_{k} + c_{j}^{(0)} \Delta t) \, \Delta t \\ &+ \sum_{j=1}^{s} \sum_{l=1}^{m} B_{i,j}^{(0)} \, \mathbf{L}^{l}(\tilde{\mathbf{x}}_{j}^{(l)}, t_{k} + c_{j}^{(1)} \Delta t) \, \Delta \hat{\beta}_{k}^{(l)}, \\ \tilde{\mathbf{x}}_{i}^{(n)} &= \hat{\mathbf{x}}(t_{k}) + \sum_{j=1}^{s} A_{i,j}^{(1)} \, \mathbf{f}(\tilde{\mathbf{x}}_{j}^{(0)}, t_{k} + c_{j}^{(0)} \Delta t) \, \Delta t \\ &+ \sum_{j=1}^{s} \sum_{l=1}^{m} B_{i,j}^{(1)} \, \mathbf{L}^{l}(\tilde{\mathbf{x}}_{j}^{(l)}, t_{k} + c_{j}^{(1)} \Delta t) \, \Delta \hat{\beta}_{k}^{(l,n)}, \\ \bar{\mathbf{x}}_{i}^{(n)} &= \hat{\mathbf{x}}(t_{k}) + \sum_{j=1}^{s} A_{i,j}^{(2)} \, \mathbf{f}(\tilde{\mathbf{x}}_{j}^{(0)}, t_{k} + c_{j}^{(0)} \Delta t) \, \Delta t \\ &+ \sum_{j=1}^{s} \sum_{l=1}^{m} B_{i,j}^{(2)} \, \mathbf{L}^{l}(\tilde{\mathbf{x}}_{j}^{(l)}, t_{k} + c_{j}^{(1)} \Delta t) \, \frac{\Delta \hat{\beta}_{k}^{(l,n)}}{\sqrt{\Delta t}}, \end{split}$$

- Again, the increments are given by the double Itô integrals.
- In the weak schemes we can use the following approximations:

$$\Delta \hat{\beta}_{k}^{(i,j)} = \begin{cases} \frac{1}{2} \left(\Delta \hat{\beta}_{k}^{(i)} \Delta \hat{\beta}_{k}^{(j)} - \sqrt{\Delta t} \, \hat{\zeta}_{k}^{(i)} \right), & \text{if } i < j, \\\\ \frac{1}{2} \left(\Delta \hat{\beta}_{k}^{(i)} \Delta \hat{\beta}_{k}^{(j)} + \sqrt{\Delta t} \, \hat{\zeta}_{k}^{(j)} \right), & \text{if } i > j, \\\\ \frac{1}{2} \left([\Delta \hat{\beta}_{k}^{(i)}]^{2} - \Delta t \right), & \text{if } i = j. \end{cases}$$

- Here only 2m 1 independent random variables are needed.
- No problems with the cross-term integrals any more.

• For example, we can choose $\Delta \hat{\beta}_{k}^{(l)}$ such that they are independent three-point distributed random variables:

$$P(\Delta \hat{\beta}_k^{(i)} = \pm \sqrt{3 \,\Delta t}) = rac{1}{6} \quad ext{and} \quad P(\Delta \hat{\beta}_k^{(i)} = 0) = rac{2}{3},$$

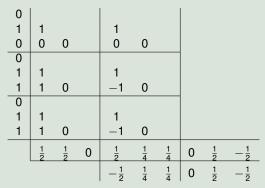
• The supporting variables $\hat{\zeta}_k^{(i)}$ such that they are independent two-point distributed random variables.

$$\mathbf{P}(\hat{\zeta}_k^{(i)} = \pm \sqrt{\Delta t}) = \frac{1}{2}.$$

Example: A weak order 2.0 method

Example (Weak order 2.0 SRK due to Rößler)

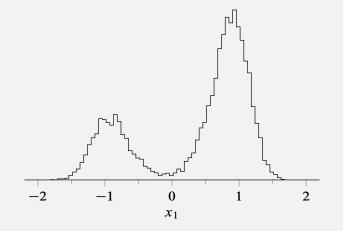
• Consider a stochastic Runge–Kutta method with the following extended Butcher tableau:



• The (rather lengthy) algorithm is written out in the lecture notes.

- We are interested in characterizing the solution at t = 20 for the initial condition of x(0) = (-3,0).
- We use the stochastic Runge–Kutta method of weak order 2.0.
- Discretization interval: $\Delta t = 2^{-4}$.
- We show the results as a histogram of $x_1(20)$ with 10,000 samples.
- With a ∆t this large, the Euler–Maruyama method does not provide plausible results.

Example: Weak SRK for Duffing van der Pol [2/2]



- Stochastic Runge–Kutta methods are derivative-free methods for solving SDEs.
- They cannot be derived as simple extensions to ordinary Runge–Kutta methods.
- You cannot get rid of the iterated Itô integral.
- The complexity of the methods grows with the approximation order.
- Higher order schemes can be practical for models with some special structure (scalar, additive, commutative, *etc.*).
- The choice between a weak and strong scheme depends on your application.